



Phase diagram of $\text{SrO-InO}_{1.5}\text{-CoO}_x$ and a new compound $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$

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ABSTRACT

$\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$, isostructural to $\text{Ca}_3\text{Co}_2\text{O}_6$, is revealed by the study of the phase relations in the system $\text{SrO-InO}_{1.5}\text{-CoO}_x$ (1000 °C). The structure of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is refined by the combination of powder X-ray and neutron diffraction. $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ crystallizes in a trigonal lattice with the cell parameters $a=b=9.59438(3)$ Å, $c=11.02172(4)$ Å with the space group $R\bar{3}c$. Its structure possesses 1D $(\text{In/Co})\text{O}_3$ chains running along the c -axis constructed by alternating face-sharing CoO_6 octahedra and $(\text{In}_{0.9}\text{Co}_{0.1})\text{O}_6$ trigonal prisms. The co-occupation of In^{3+} and Co^{3+} at the trigonal prismatic site is evidenced by elementary analysis and determined by the structure refinement. $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is paramagnetic, and the susceptibility is consistent with the occupation of Co^{3+} at 10% of the trigonal prismatic positions in a high spin state (HS, $S=2$). The HS Co^{3+} is well separated by diamagnetic CoO_6 octahedra and InO_6 trigonal prisms and shows a g factor of 2.0 in the magnetic measurements.

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1. Introduction

Hexagonal perovskite related compounds with the general formula $A_3B'\text{BO}_6$ have attracted considerable attention due to their interesting structural features and physical properties [1–8]. In this family, A is typically Ca or Sr, and B' and B can be the same, as in $\text{Ca}_3\text{Co}_2\text{O}_6$, or different cations, with $B'=Y, \text{Sc}, \text{In}, \text{Na}, \text{Mg}, \text{Ni}, \text{Cd}, \text{Zn}, \text{Co}, \text{Sm}, \text{Eu}, \text{Tb}, \text{Dy}, \text{Ho}, \text{Er}, \text{Tm}, \text{Yb}, \text{Lu}, \text{Pb}$, and $B=\text{Ru}, \text{Rh}, \text{Ir}, \text{Pd}, \text{Pt}, \text{Cr}, \text{Ni}, \text{Bi}$ [6]. The structure of $A_3B'\text{BO}_6$ consists of infinite chains formed by alternating face-sharing $B'\text{O}_6$ trigonal prisms and BO_6 octahedra and each chain is surrounded by six chains separated by the A cations. These compounds possess quasi-1D structures in which the interaction within the chains is much stronger than that between the chains. Among them, $\text{Ca}_3\text{Co}_2\text{O}_6$ is a well investigated one [1,3,9–14]. It presents a ferromagnetic ordering of the magnetic Co^{3+} cations along the chains and antiferromagnetic correlations between chains at low temperature. Its magnetization curve displays steps at low temperature, which is considered as the result of interchain frustration. Due to the similar chemistry of Ca^{2+} and Sr^{2+} the larger size of Sr^{2+} accommodates more B and/or B' cations, so substitution of Ca^{2+} by Sr^{2+} results in a series of $\text{Sr}_3B'\text{BO}_6$ compounds [6].

In an effort to explore new compounds in the $\text{Sr}_3B'\text{BO}_6$ family, we focused on the system $\text{SrO-InO}_{1.5}\text{-CoO}$. In the related systems, no binary compound was revealed in the system $\text{InO}_{1.5}\text{-CoO}$, and only one compound, SrIn_2O_4 , with a CaFe_2O_4 -type structure was reported

in $\text{SrO-InO}_{1.5}$ [15], while the system SrO-CoO_x was complicated [5,7,16–22], in which $\text{SrCoO}_{2.52}$, $\text{Sr}_6\text{Co}_5\text{O}_{15}$, $\text{Sr}_5\text{Co}_4\text{O}_{12}$, $\text{Sr}_{24}\text{Co}_{19}\text{O}_{57}$, $\text{Sr}_4\text{Co}_3\text{O}_9$, $\text{Sr}_{14}\text{Co}_{11}\text{O}_{33}$, and $\text{Sr}_3\text{Co}_2\text{O}_{7-y}$ ($0.94 \leq y \leq 1.22$) were revealed. $\text{SrCoO}_{2.52}$ (JCPDF 40-1018) has an oxygen deficient hexagonal perovskite structure. $\text{Sr}_6\text{Co}_5\text{O}_{15}$, $\text{Sr}_5\text{Co}_4\text{O}_{12}$, $\text{Sr}_{24}\text{Co}_{19}\text{O}_{57}$, $\text{Sr}_4\text{Co}_3\text{O}_9$, and $\text{Sr}_{14}\text{Co}_{11}\text{O}_{33}$ belong to the hexagonal perovskite family Sr_xCoO_3 [6–7,18–22] and $\text{Sr}_3\text{Co}_2\text{O}_{7-y}$ ($0.94 \leq y \leq 1.22$) is an oxygen-deficient $n=2$ member of the Ruddlesden–Popper phase [23]. These strontium–cobalt complex oxides present various coordination configurations of cobalt, including octahedra, trigonal prisms, and intermediate polyhedra formed from octahedra and trigonal prisms, and accordingly there are a variety of magnetic and other physical properties [5,7,16–23].

This work presents the phase diagram of $\text{SrO-InO}_{1.5}\text{-CoO}_x$ at 1000 °C. A new ternary compound $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$, isostructural to $\text{Ca}_3\text{Co}_2\text{O}_6$, is revealed. In the 1D $B'\text{BO}_6$ chain formed by face-sharing octahedra and trigonal prisms, octahedral positions are fully occupied by Co^{3+} in the low spin state (LS, $S=0$) and trigonal prismatic positions are occupied by In^{3+} and Co^{3+} in the ratio of 9:1, where Co^{3+} are in the high spin state (HS, $S=2$) and well separated by In^{3+} and Co^{3+} (LS). This compound is paramagnetic due to the 10% occupancy of Co^{3+} (HS) in the trigonal prismatic positions.

2. Experimental

2.1. Materials and preparation

The samples were synthesized by conventional solid state reaction. SrCO_3 , In_2O_3 , and Co_3O_4 of analytical grade were mixed

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From the neutron diffraction data, the positions and the occupancies of oxygen atoms can be accurately determined. The structure refinement indicates that In^{3+} and Co^{3+} are located at the center of the trigonal prisms randomly in a ratio of 0.92(1):0.08(1) and all the octahedral positions are occupied by Co^{3+} atoms, which is consistent with the elementary analysis result $n(\text{In}):n(\text{Co})=0.91:1.09$. It is noteworthy that we did not

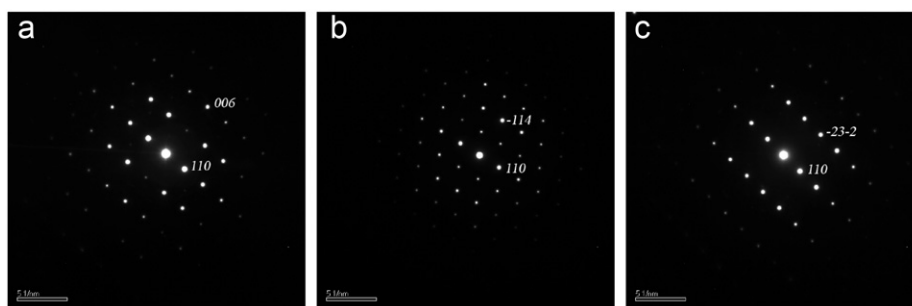


Fig. 2. Electron diffraction patterns of a tilt series of one crystallite for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ along the zone axes: (a) $[1 - 10]$; (b) $[2 - 21]$; (c) $[-225]$.

Table 1

Neutron data collection conditions, crystallographic data and results of Rietveld analysis for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ (298 K).

Crystal data	
Chemical formula	$\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$
M_r (g mol ⁻¹)	527.02
Space group (number)	$R\bar{3}c$ (167)
Lattice parameters (Å, deg.)	$a=9.59438(3)$, $c=11.02172(4)$
V (Å ³)	878.65(1)
Z	6
D_x (mg m ⁻³)	5.97(1)
Radiation type, λ (Å)	Neutron (unpolarized), 1.494
μR	0.3
Data collection	
Diffractometer	High resolution powder diffractometer for thermal neutrons
Data collection method	Debye–Scherrer geometry
Specimen mounting	Vanadium container
Detector	PSD
Data collection mode	Transmission
2θ min., max., step (°)	5, 164.9, 0.05
Refinement	
R_p , R_{wp} , χ^2	0.030, 0.039, 4.25
R_{expected}	0.019
Excluded regions (°)	5.00–11.00
$N - P + C$	3047
Computer program	Fullprof 2009

Table 2

Atomic coordinates, occupancies and isotropic thermal displacement parameters for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ (298 K).

Atom	Wyckoff positions	x	y	z	Occupancy	B_{eq} (Å ²)
Sr	18e	0.3690(1)	0.00000	0.25000	1	0.49(1)
In	6a	0.00000	0.00000	0.25000	0.92(1)	0.49(4)
Co2	6a	0.00000	0.00000	0.25000	0.08(1)	0.49(4)
Co1	6b	0.00000	0.00000	0.00000	1	0.51(5)
O	36f	0.1713(1)	0.0211(1)	0.1100(1)	1	0.54(1)

observe any evidence of the simple stoichiometric compound “ $\text{Sr}_3\text{InCoO}_6$ ”, which reflects that the formation of the structure relies strongly on the radii of the $B(B')$ -site cations.

The Co–O bond lengths and O–Co–O bond angles in the CoO_6 octahedra are consistent with the data previously reported for $\text{Ca}_3\text{Co}_2\text{O}_6$ [10]. Since the octahedron is elongated along the c direction, the Co^{3+} cation in this position has anisotropic thermal factors with $U_{33}=0.014 \text{ Å}^2$, $U_{11}=U_{22}=0.003 \text{ Å}^2$ (See Table S3 and Fig. S5). The extrapolation of the thermal factors to 0 K shows that U_{11} and U_{22} tend to zero, but U_{33}

Table 3

Selected bond distances (Å) and bond angles (deg.) for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ (298 K).

CoO₆ octahedra	
Co–O	1.970(1) × 6
O–Co–O	93.92(4) × 6; 86.08(3) × 6
In(Co)O₆ trigonal prism	
In(Co)–O	2.189(1) × 6
O–In(Co)–O	90.37(4) × 3; 75.80(3) × 6 130.12(4) × 3; 147.10(5) × 3

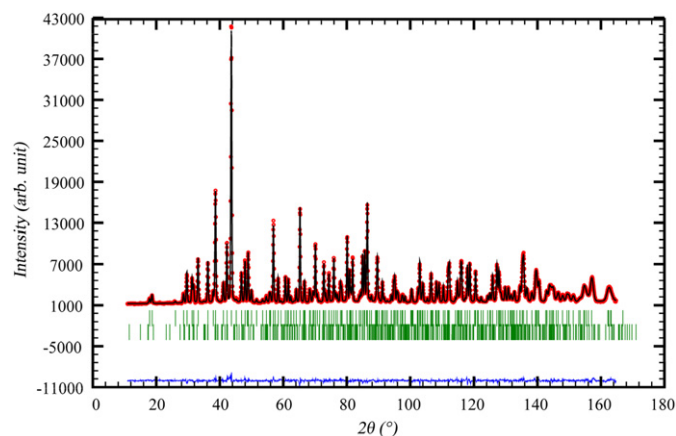


Fig. 3. Rietveld refinement plot of the powder neutron diffraction profile of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ at 298 K. The circles are the observed data and the solid line is for the calculated pattern, the marks below the diffraction pattern are reflection positions of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ (upper) and SrIn_2O_4 (lower, ~3 wt%), and the difference curve is shown at the bottom.

has a residual value of $\sim 0.01 \text{ Å}^2$, which is generally considered to show the existence of a static displacive disorder in the structure [26]. In this case, it relates to the elongated extension of the CoO_6 octahedron along the c direction. For the InO_6 trigonal prism, the bond lengths and angles are reasonable. However, considering the occupation of Co^{3+} in trigonal prismatic positions, the space is quite large for a Co^{3+} cation, enabling the Co^{3+} to take the $S=2$ high spin state.

In the $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ structure (Fig. 4), the chains are formed by alternating face-sharing MO_6 ($M=\text{In}^{3+}$, Co^{3+}) trigonal prisms and CoO_6 octahedra. Each chain is surrounded by six chains, and neighboring chains are shifted with respect to each other by a $1/3^*c$ (translation) along the c direction. The $B'\text{O}_6$ trigonal prism can be considered as the result of the substitution of two B and three interstitial oxygen atoms in the ABO_3 hexagonal perovskite by one B' . In such a structure, the height of the $B'\text{O}_6$ trigonal prism is

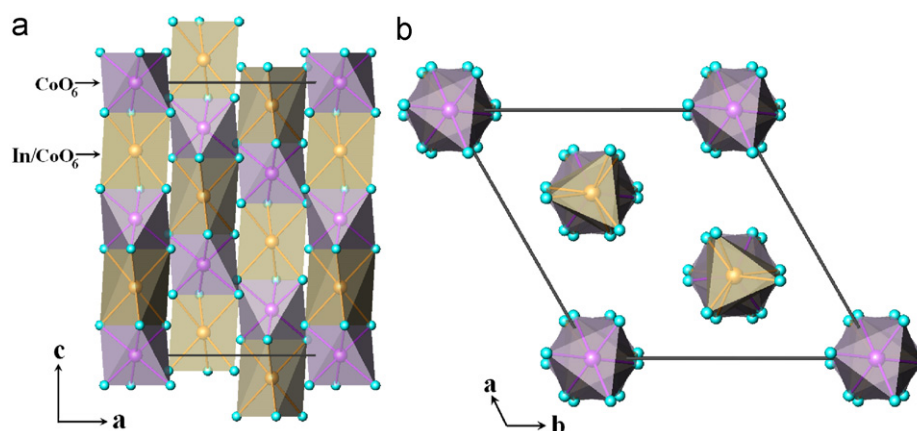


Fig. 4. Distribution of 1D chains in the structure of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ viewed along the directions $[010]$ (a) and $[001]$ (b).

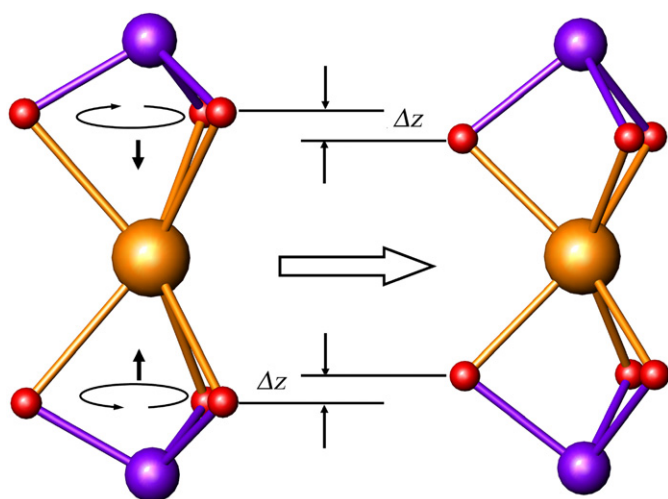


Fig. 5. Schematic pattern of the shift and rotation of oxygen atoms. The small red balls stand for oxide anions, the large and medium balls stand for B' (in trigonal prism) and B (in octahedron) cations, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

roughly twice that of BO_6 octahedron. In real structures the oxygen atoms of $B'\text{O}_6$ rotate and shift with respect to B' , and the shift distance and the rotation angle describe the deviation from the ideal coordination. Here the relative shift $\delta_z (= \Delta z/c)$, the quotient of the shift distance Δz (shown in Fig. 5) divided by c parameter, is used to compare the distortion of the chain. We summarized the structure data of some $\text{Sr}_3B'\text{BO}_6$ compounds [6,27–31] and found that δ_z had an approximately linear relationship with the ratio $r_{\text{TP}}/r_{\text{OC}}$; here r_{TP} and r_{OC} stand for the radius of B' and B cations, as shown in Fig. 6. The shift depends on the radii of B' and B atoms. The closer the B' and B atoms are, the more obvious the shift is. In $\text{Sr}_3B'\text{BO}_6$ [6,27–31], δ_z is in the range of 0.02–0.04. The distortion of the octahedron and trigonal prisms for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is comparable to the others.

3.3. Magnetic properties of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$

In $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$, only 1/10 of trigonal prismatic positions are occupied by Co^{3+} (HS); then the formula weight containing one mole of Co^{3+} (HS) is ten times of that of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ and the molar susceptibility of the Co^{3+} (HS) is calculated accordingly. The dependence of magnetic susceptibility on temperature is shown

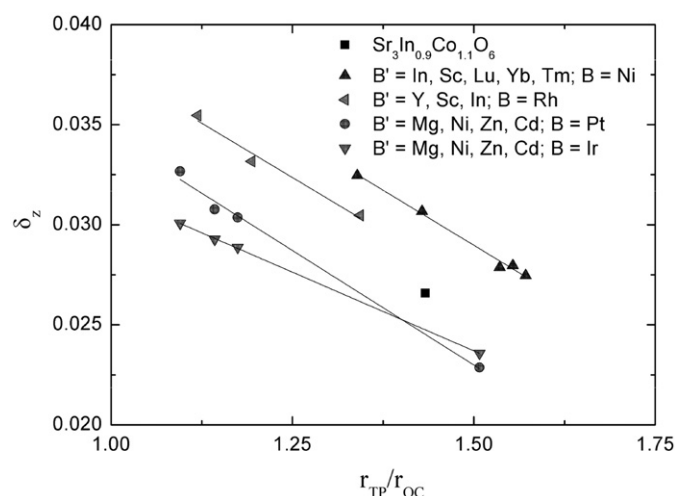


Fig. 6. Relative shift ($\delta_z = \Delta z/c$) in $\text{Sr}_3B'\text{BO}_6$.

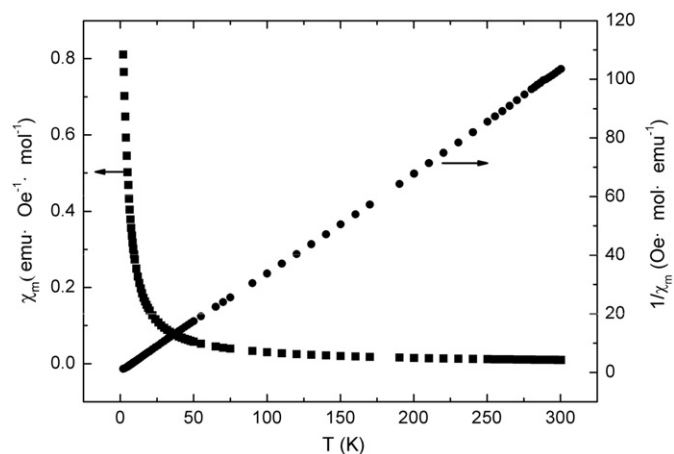


Fig. 7. Plots of magnetic susceptibility and its inverse with temperature for $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$.

in Fig. 7. The plot of the inverse magnetic susceptibility to temperature fits Curie's Law very well, indicating that $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is paramagnetic. The neutron diffraction data of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ at low temperatures also confirm the paramagnetic behavior. The experiments carried out down to 1.5 K did not show any signs of

long-range magnetic order. The Curie constant C is $2.9 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$. According to the equation $8C = g^2 S(S+1)$, we can deduce that the g factor of $\text{Co}^{3+}(\text{HS})$ in $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is about 2.0.

The magnetism of $\text{Co}^{3+}(\text{HS})$ in $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is isolated instead of being coupled. Its nearest neighbor in the chain is $\text{Co}^{3+}(\text{LS})$ and the next nearest neighbor is In^{3+} with 90% possibility and $\text{Co}^{3+}(\text{HS})$ with 10% possibility. The increase of the content of $\text{Co}^{3+}(\text{HS})$ in trigonal prisms is helpful to the formation of segments of $\text{Co}^{3+}(\text{HS})$ – $\text{Co}^{3+}(\text{LS})$ – $\text{Co}^{3+}(\text{HS})$, and as a consequence enhances the interaction between $\text{Co}^{3+}(\text{HS})$ cations. Since in the system SrO – $\text{InO}_{1.5}$ – CoO_x at 1000°C , $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ appears as a sharp compound, the content of Co^{3+} may be adjusted by the simultaneous substitution of Sr^{2+} by other cations.

4. Conclusions

In the investigation of the phase relations of the system SrO – $\text{InO}_{1.5}$ – CoO at 1000°C in atmosphere, a new ternary compound $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$, isostructural to $\text{Ca}_3\text{Co}_2\text{O}_6$, is revealed. The structure of $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ is confirmed by Rietveld analysis of powder X-ray and neutron diffraction data. $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$ crystallizes in a trigonal system with the cell parameters $a=b=9.59438(3) \text{ \AA}$, $c=11.02172(4) \text{ \AA}$ in the space group $R\bar{3}c$. Its structure contains a 1-D chain constructed by alternating face-sharing CoO_6 octahedra and $(\text{In}_{0.9}\text{Co}_{0.1})\text{O}_6$ trigonal prisms. The co-occupation of In^{3+} and Co^{3+} in the trigonal prism is evidenced by elementary analysis and determined by the refinement based on neutron diffraction data. In $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$, the $\text{Co}^{3+}(\text{HS})$ in trigonal prismatic positions are well separated by diamagnetic Co^{3+}O_6 octahedra and InO_6 trigonal prisms, such that the interaction between $\text{Co}^{3+}(\text{HS})$ is very weak, and no magnetic ordering was observed even down to 1.5 K . The magnetic measurements indicate that the compound is paramagnetic, and the susceptibility is consistent with 10% of the trigonal prismatic positions being occupied by Co^{3+} cations with the high spin $S=2$. Magnetic data give the value of the g factor of 2.0 for $\text{Co}^{3+}(\text{HS})$ in $\text{Sr}_3\text{In}_{0.9}\text{Co}_{1.1}\text{O}_6$.

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Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.jssc.2011.02.024.

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